High-Dimensional SQKD

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New high-dimensional semi-quantum key distribution protocol $^{\rm 1}$

Simpler method for security analysis

Proof of information theoretic security

¹https://arxiv.org/abs/1907.11340



Quantum Key Distribution(QKD)



A Concrete Example





QKD: All parties have advanced quantum capabilities. What if Bob can't measure in the *X* basis? Semi-Quantum Key Distribution (Boyer et al., 2007)

Bridge the gap between Classical and Quantum realms

Use less expensive hardware

Fallback option for fully fledged QKD



But It's Too...



A Concrete Example



SQKD

QKD with High-dimensional(HD) systems

Naturally carries more information

More robust against quantum cloning

More noise resistant



Earlier works on HD-QKD :

High-dimensional quantum key distribution based on mutually partially unbiased bases (Wang et al., 2020)

Provably secure and high-rate quantum key distribution with time-bin qudits (Islam et al., 2017)

Security proof for quantum key distribution using qudit systems (Sheridan et al., 2010)



Can we use HD-systems in SQKD scenario and still prove information theoretic security?

A key is secure if Alice's and Bob's keys are the same and Eve has no knowledge about it.

Let a joint state among Alice, Bob and Eve be:

$$\rho_{XYE}^{real} = \sum_{k_A, k_B \in \{0,1\}^l} \Pr(k_A, k_B) |k_A \rangle \langle k_A| \otimes |k_B \rangle \langle k_B| \otimes \rho_E^{(k_A, k_B)},$$

and another desired state be:

$$\rho_{XYE}^{ideal} = \frac{1}{2^{l}} \sum_{k \in \{0,1\}^{m}} |k\rangle \langle k|_{A} \otimes |k\rangle \langle k|_{B} \otimes \rho_{E}.$$

Then the final key k is said to be ϵ -secure if (Renner 2005)

$$\frac{1}{2}||\rho_{XYE}^{\textit{real}} - \rho_{XYE}^{\textit{ideal}}||_1 \le \epsilon,$$

where $||A||_1 := Tr(\sqrt{A^{\dagger}A})$ is the trace norm of A.

 $\epsilon = \epsilon' + \epsilon''$ -security also implies that it is:

$$\epsilon'$$
-correct := $\Pr(k_A \neq k_B) \leq \epsilon'$,

and if C is communication transcript in IR,

$$\epsilon''$$
-secret := $\frac{1}{2} ||\rho_{XCE}^{real} - \rho_{XCE}^{ideal}||_1 \le \epsilon''$

Renner (2005) proved that:

$$||\rho_{XCE}^{real} - \rho_{XCE}^{ideal}||_1 \leq 2^{-\frac{1}{2}(H_{min}(X|CE) - I)},$$

where / is the final length of the key. We want the r.h.s to be at most $2\epsilon''$.

Solving for *I*, we get:

$$I \leq H_{min}(X|CE) + 2\log(2\epsilon'')$$

Let's take the equality to get the maximum length:

$$I = H_{min}(X|CE) + 2\log(2\epsilon'')$$

$$\geq H_{min}(XC|E) - H_{max}(C) + 2\log(2\epsilon'') \text{ [Chain rule]}$$

$$\geq H_{min}(X|E) + H_{min}(C|X) - H_{max}(C) + 2\log(2\epsilon'')$$

$$\geq H_{min}(X|E) - (H_{max}(C) - H_{min}(C|X)) + 2\log(2\epsilon'')$$

Now, as we are interested in the asymptotic scenario, where:

$$\frac{1}{n}H_{min}(X|E) = H(X|E), \text{ and } \frac{1}{n}(H_{max}(C) - H_{min}(C|X)) = H(X|Y).$$

So, finally:

$$I \ge H(X|E) - H(X|Y) + 2\log(2\epsilon'')$$

Finally, if ρ is one key-iteration of a protocol, the key-rate is defined as:

key-rate :=
$$\frac{I}{n} = min(H(X|E)_{\rho} - H(X|Y)_{\rho})$$



The Protocol





SQKD: $\left|0\right\rangle,\left|1\right\rangle,\left|+\right\rangle,\left|-\right\rangle$ are two dimensional.

HD-SQKD: $|z\rangle \in \{|0\rangle, |1\rangle \dots |N-1\rangle\},\ |x\rangle \in \mathcal{F}\{|0\rangle, |1\rangle \dots |N-1\rangle\}$ where \mathcal{F} is Quantum Fourier transformation.

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Problem: Two-way SQKD analysis and density matrix computation is too complex





Solution: Reduction to One Way Protocol



Two-Way SQKD



One-Way SQKD

Theorem

Let (U_F, U_R) be a collective attack against HD-SQKD. Then, there is an attack against the OW-SQKD protocol such that, Eve gets no advantage in either scenario.

HD-SQKD	OW-SQKD
1. A prepares $ z\rangle$ or $ x\rangle$, sends	1. Bob prepares and sends $ \phi_R angle$
to Bob	or $ \phi_{M\!R} angle$ if he wants to reflect or
	measure respectively
2. Eve attacks with U_F	2. Eve attacks with U
3. Bob measures or resends in	3. Alice measures A_1 and A_2
\mathcal{Z} basis	registers in $\mathcal Z$ or $\mathcal X$ basis
4. Eve attacks with U_R	
5. Alice measures the returning	
n qubits in the preparation basis	

$$|z\rangle \in \{|0\rangle, |1\rangle, \dots, |N-1\rangle\}, |x\rangle = \mathcal{F}|z\rangle$$
(1)

$$|\phi_R\rangle = \sum_{b=0}^{2^n-1} \sqrt{\rho(b)} |b,b\rangle_{A_1A_2} \otimes |0\rangle_B$$
(2)

$$|\phi_{MR}\rangle = \sum_{b=0}^{2^n-1} \sqrt{\rho(b)} |b, b, b\rangle_{A_1A_2B}$$
(3)

Proof











In two-way case, let Alice's choices are $|0\rangle\,,|1\rangle\,,|2\rangle\,,|3\rangle$ and she chooses $|1\rangle$ to send to Bob. Eve attacks then:

$$U_{F} \left| 1
ight
angle = \left| 0, e_{10}
ight
angle + \left| 1, e_{11}
ight
angle + \left| 2, e_{12}
ight
angle + \left| 3, e_{13}
ight
angle$$

Bob measures and finds a $|2\rangle$ with probability $\langle e_{12}|e_{12}\rangle$. Then, one-way case, R_w must recreate all the scenarios where Bob could measure a $|2\rangle$. Specifically,

$$\begin{array}{l} |0\rangle \rightarrow |2\rangle \,, \, \text{with probability} \, \left\langle e_{02}|e_{02}\right\rangle \\ |1\rangle \rightarrow |2\rangle \,, \, \text{with probability} \, \left\langle e_{12}|e_{12}\right\rangle \\ |2\rangle \rightarrow |2\rangle \,, \, \text{with probability} \, \left\langle e_{22}|e_{22}\right\rangle \\ |3\rangle \rightarrow |2\rangle \,, \, \text{with probability} \, \left\langle e_{32}|e_{32}\right\rangle \end{array}$$

So,

$$egin{aligned} R_{w} \left| 2,2
ight
angle &= rac{\left| 0,2,e_{02}
ight
angle + \left| 1,2,e_{12}
ight
angle + \left| 2,2,e_{22}
ight
angle + \left| 3,2,e_{32}
ight
angle }{\sqrt{4 \cdot
ho(2)}} \end{aligned}$$

Key-rate Computation



Only Measure/Resend (M/R) rounds are key-generating-rounds. Reflect rounds are used for noise estimation. So the goal is to upper bound Eve's uncertainty about Alice's register in M/R case.



But it's not observable!



Eve's uncertainty about the reflect case is not observable either.



But Alice can measure the uncertainty in the reflect case!



Entropic uncertainty relation (Berta et al., 2009): For any density operator $\rho_{A_1A_2E}$ and two measurements \mathcal{Z} and \mathcal{F} ,

$$H(A_1^Z|E) + H(A_1^F|A_2) \ge n$$



We can bound Eve's uncertainty in Reflect case



Continuity bound(Winter, 2015): For states ρ and μ on a Hilbert space $A \otimes E$, if $\frac{1}{2} ||\rho - \sigma|| \le \delta \le 1$ then

$$|H(A_1^Z|E)_{\rho} - H(A_1^Z|E)_{\mu}| \le 2\delta \log |A_1^Z| + (1+\delta)h(\frac{\delta}{1+\delta})$$



We can bound Eve's uncertainty in M/R case



In our case, δ is linear function.

 $\delta = f(\text{noise, eigenvalues, dimension}),$

and key-rate r is:

$$r \geq n(1-\delta) - (1+\delta)H(rac{\delta}{1+\delta}) - 2Q\log_2(2^n-1) - 2H(Q),$$

where, n is the number of qubits sent, δ is the trace distance, Q is the noise parameter.



Key-rate of our HD-SQKD protocol



I'VE DISCOVERED A WAY TO GET COMPUTER SCIENTISTS TO LISTEN TO ANY BORING STORY.

